

Exterior energy bounds for wave equations and applications

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1 Energy-critical wave equation

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The equation

$$(NLW) \quad \begin{cases} \partial_t^2 u - \Delta u = |u|^{\frac{4}{N-2}} u \\ \vec{u}|_{t=0} = (u_0, u_1) \in \mathcal{H} = \dot{H}^1(\mathbb{R}^N) \times L^2(\mathbb{R}^N) \end{cases}$$

where $u : [0, T[\times \mathbb{R}^N \rightarrow \mathbb{R}$, $N \geq 3$.

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The energy

$$E(\vec{u}) = \frac{1}{2} \int_{\mathbb{R}^N} |\nabla_{t,x} u(t)|^2 - \frac{N-2}{2N} \int_{\mathbb{R}^N} |u(t)|^{\frac{2N}{N-2}}$$

is conserved, where $|\nabla_{t,x} u|^2 = (\partial_t u)^2 + \sum_{j=1}^N (\partial_{x_j} u)^2$.

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Transformations: if u is a solution and $\lambda > 0$, $t_0 \in \mathbb{R}$, $\varepsilon \in \{\pm 1\}$, so is

$$v(t, x) = \frac{\varepsilon}{\lambda^{N/2-1}} u\left(\frac{t-t_0}{\lambda}, \frac{x-x_0}{\lambda}\right).$$

$$\|\vec{v}(t_0)\|_{\mathcal{H}} = \|\vec{u}(0)\|_{\mathcal{H}}, \quad E(\vec{v}) = E(\vec{u}).$$

Travelling waves

Stationary solutions of (NLW):

$$(E) \quad -\Delta Q = |Q|^{\frac{4}{N-2}} Q, \quad Q: \mathbb{R}^N \rightarrow \mathbb{R}, \quad Q \in \dot{H}^1(\mathbb{R}^N).$$

Minimal energy nonzero solution of (E) (**ground state**):

$$W(x) = \left(1 + \frac{|x|^2}{N(N-2)} \right)^{1-\frac{N}{2}}$$

(unique radial solution up to transformations).

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Solitary waves ($\mathbf{p} \in \mathbb{R}^N$, $|\mathbf{p}| < 1$): $Q_{\mathbf{p}}(t, x) = Q_{\mathbf{p}}(0, x - t\mathbf{p})$, where

$$Q_{\mathbf{p}}(0, x) = Q \left(\left(\frac{1}{\sqrt{1-p^2}} - 1 \right) \frac{(\mathbf{p} \cdot x)\mathbf{p}}{p^2} + x \right)$$

Type II blow-up solutions

Definition. u is a *type II blow-up solution* of (NLW) when the maximal time of existence T_+ of u is finite and

$$\limsup_{t \rightarrow T_+} \|\vec{u}(t)\|_{\mathcal{H}} < \infty.$$

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Existence of solutions of the form:

$$\vec{u}(t) = \left(\frac{1}{\lambda(t)^{\frac{N-2}{2}}} W\left(\frac{\cdot}{\lambda(t)}\right), 0 \right) + (v_0, v_1), \quad t \rightarrow T_+,$$

where $(v_0, v_1) \in \mathcal{H}$ et $\lambda(t) \ll T_+ - t$, [Krieger Schlag Tataru 09, 14] ($N = 3$) and also: [Hillairet Raphaël 12] ($N = 4$), [Jendrej 2015].

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Open question: Type II solutions with more than one bubble, or other bubbles than W . See: [Jendrej], [Martel, Merle] (global case) and also [Côte, Zaag] (1D), [Côte, Martel] (KG).

Soliton resolution conjecture for type II blow-up

Conjecture. Let u be a *radial*, type II blow-up solution of (NLW). Then there exists $J \geq 1$ and:

- $(v_0, v_1) \in \mathcal{H}$,
- signs $\varepsilon_j \in \{\pm 1\}$, $j = 1 \dots J$,
- parameters $\lambda_j(t)$, $0 < \lambda_1(t) \ll \lambda_2(t) \ll \dots \ll \lambda_J(t) \ll T_+ - t$,

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$$\vec{u}(t) = (v_0, v_1) + \sum_{j=1}^J \left(\frac{\varepsilon_j}{\lambda_j^{\frac{N}{2}-1}(t)} W \left(\frac{x}{\lambda_j(t)} \right), 0 \right) + \vec{w}(t),$$

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In the *nonradial case*: possibly several blow-up points, more general solitons (replace $\pm W$ by $Q_{\mathbf{p}_j}^j$), space translations allowed.

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Asymptotic behaviour for linear waves

Consider:

$$(LW) \quad \begin{cases} \partial_t^2 u_L - \Delta u_L = 0, & x \in \mathbb{R}^N, t \in \mathbb{R} \\ \vec{u}_L|_{t=0} = (u_0, u_1) \in \mathcal{H} = \dot{H}^1(\mathbb{R}^N) \times L^2(\mathbb{R}^N). \end{cases}$$

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Then (see e.g. [Friedlander]) there exist $G_{\pm} \in L^2(\mathbb{R} \times S^{N-1})$ such that:

$$\begin{aligned} \lim_{t \rightarrow +\infty} \int_0^{+\infty} \int_{S^{N-1}} \left| r^{\frac{N-1}{2}} \partial_r u_L(t, r\omega) \mp G_{\pm}(r-t, \omega) \right|^2 \\ + \left| r^{\frac{N-1}{2}} \partial_t u_L(t, r\omega) + G_{\pm}(r-t, \omega) \right|^2 dr d\omega = 0. \end{aligned}$$

Furthermore (denoting ∂ the tangential derivative):

$$\lim_{t \rightarrow +\infty} \int \frac{1}{|x|^2} |u_L(t, x)|^2 + |\partial u_L(t, x)|^2 + |u_L(t, x)|^{\frac{2N}{N-2}} dx = 0$$

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$(u_0, u_1) \mapsto G_+$ and $(u_0, u_1) \mapsto G_-$ are **isometries** between \mathcal{H} and $L^2(\mathbb{R} \times S^{N-1})$.

Equirepartition of the energy

Theorem [TD, Kenig, Merle 2012]. Assume N is odd. Let u_L be the solution of the linear wave equation with data (u_0, u_1) . Then

$$\sum_{\pm} \lim_{t \rightarrow \pm\infty} \int_{\{|x| \geq |t|\}} |\nabla_{t,x} u_L(t, x)|^2 dx = \int_{\mathbb{R}^N} |\nabla_{t,x} u_L(0, x)|^2 dx.$$

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More precisely:

- if $u_1 = 0$ then $G_+(\eta, \omega) = (-1)^{\frac{N+1}{2}} G_+(-\eta, -\omega)$.
- if $u_0 = 0$ then $G_+(\eta, \omega) = (-1)^{\frac{N-1}{2}} G_+(-\eta, -\omega)$,

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The lower bound (\geq) does not hold in even dimension ([Côte, Kenig, Schlag 2014]), even if a multiplicative constant is allowed.

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Nonradiative solution

Definition. A (global) solution u of (NLW) is *nonradiative* when

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Solitary waves are nonradiative.

We look for *rigidity theorems* of the form: u nonradiative $\implies u$ is a solitary wave.

This is difficult: it implies in particular the nonexistence of pure multisolitons (typically false for integrable equations). See [Martel Merle 2017].

Rigidity theorem for small data

Theorem [TD, Kenig, Merle 2012]. Assume N is odd. Then there exists $\varepsilon_0 > 0$ such that if u is a solution of (NLW) with:

$$\|(u_0, u_1)\|_{\mathcal{H}} < \varepsilon_0,$$

then

$$\sum_{\pm} \lim_{t \rightarrow \pm\infty} \int_{\{|x| \geq |t|\}} |\nabla_{t,x} u(t, x)|^2 dx \geq \frac{1}{2} \int_{\mathbb{R}^N} |\nabla_{t,x} u(0, x)|^2 dx.$$

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Corollary [TD, Kenig, Merle 2012] Soliton resolution for type II blow-up solutions holds with the additional assumption:

$$\limsup_{t \rightarrow T_+} \|\vec{u}(t)\|_{\mathcal{H}}^2 \leq \|W\|_{\dot{H}^1}^2 + \varepsilon^2.$$

See works of [Krieger, Nakanishi, Schlag] for a complete description of the dynamics close to the ground state.

Improved exterior energy bound for radial solution

Let $R > 0$. Denote by \mathcal{H}_R the space of radial functions in $(\dot{H}^1 \times L^2)(\{r > R\})$, and π_R^\perp the orthogonal projection, in \mathcal{H}_R , on the orthogonal of $P_R = \text{span}((1/r, 0))$.

Proposition. *Let u_L be a radial solution of the linear wave equation, in space dimension 3, with initial data (u_0, u_1) . Then*

$$\sum_{\pm} \lim_{t \rightarrow \pm\infty} \int_{R+|t|}^{+\infty} (\partial_{t,r} u_L(t, r))^2 r^2 dr = \|\pi_R^\perp(u_0, u_1)\|_{\mathcal{H}_R}^2.$$

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Generalization to [other odd dimension](#), with P_R defined as above: [\[Kenig, Lawrie, Baoping Liu, Schlag 2015\]](#). Note that:

$$\dim P_R = \frac{N-1}{2}$$

Rigidity theorem

Theorem. *Assume $N = 3$. If u is a radial non radiative solution, then $u = 0$ or there exist $\lambda > 0$, $\iota \in \{\pm 1\}$ such that $u(t, x) = \frac{\iota}{\lambda^{1/2}} W\left(\frac{x}{\lambda}\right)$.*

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$$\frac{\iota}{\lambda^{1/2}} W\left(\frac{x}{\lambda}\right) \sim \frac{\sqrt{3}\lambda^{1/2}}{|x|}, \quad |x| \rightarrow \infty.$$

First step of the proof: $\exists \ell \in \mathbb{R}$ such that

$$\lim_{r \rightarrow \infty} ru_0(r) = \ell.$$

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Difficulty in higher odd dimensions due to the higher dimension of P_R !

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Lower bound for well-prepared initial data

Lemma. *Let $\gamma \in (0, 1)$. There exists $\varepsilon = \varepsilon(\gamma) > 0$ with the following property. Let u_L be a solution of (LW) with initial data (u_0, u_1) such that*

$$\begin{cases} (u_0, u_1) \in \dot{H}^1 \times L^2 & \text{if } N \geq 3 \\ |\nabla u_0| \in L^2, u_1 \in L^2 \text{ and } u_0(x) = u_\infty \text{ for large } |x| & \text{if } N = 2 \end{cases}$$

(where $u_\infty \in \mathbb{R}$) and:

$$\|(\nabla u_0, u_1)\|_{L^2(B_{1+\varepsilon}^c \cup B_{1-\varepsilon})} + \|\partial_t u_0\|_{L^2} + \|\partial_r u_0 + u_1\|_{L^2} \leq \varepsilon \|(\nabla u_0, u_1)\|_{L^2}.$$

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(*outgoing condition*). Then, for all $t \geq 0$,

$$\int_{|x| \geq \gamma+t} |\nabla_{x,t} u_L|^2(x, t) dx \geq \gamma \|(\nabla u_0, u_1)\|_{L^2}^2.$$

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Application: soliton resolution up to a sequence of times for (NLW) [TD, Jia, Kenig, Merle 2017]. Soliton resolution for small type II blow-up solutions for wave maps. [TD, Jia, Kenig, Merle 2016], using [Grinis 2016].

Wave maps

$$(WM) \quad \begin{cases} \partial_t^2 u - \Delta u = (|\nabla u|^2 - |\partial_t u|^2) u, & x \in \mathbb{R}^2 \\ \vec{u}|_{t=0} = (u_0, u_1), & u_0 \cdot u_1 = 0, \quad |u_0| = 1. \end{cases}$$
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Blow-up solutions [Krieger, Schlag and Tataru 2008], [Raphaël, Rodnianski 2012], [Jendrej 2016].

Well prepared initial data for wave maps

Theorem. *Let $\gamma \in (0, 1)$. Then there exists $\varepsilon = \varepsilon(\gamma) > 0$ with the following property. Let u be a classical solution of (WM) with initial data (u_0, u_1) such that*

$$E_M(u_0, u_1) \leq \varepsilon$$

and

$$\|(\nabla u_0, u_1)\|_{L^2(B_{1+\varepsilon}^c \cup B_{1-\varepsilon})} + \|\not\partial u_0\|_{L^2} + \|\partial_r u_0 + u_1\|_{L^2} \leq \varepsilon \|(\nabla u_0, u_1)\|_{L^2}.$$

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Then for all $t \geq 0$,

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Small type II blow-up solutions for wave maps

Theorem. *Let u be a classical solution of (WM) such that $E_M(\vec{u}(0)) < E_M(S, 0) + \epsilon_0^2$, blowing up in finite time T_+ at $x = 0$. Then there exists $\mathbf{p} \in \mathbb{R}^2$ such that $|\mathbf{p}| \ll 1$, $x(t) \in \mathbb{R}^2$, $\lambda(t) > 0$ with*

$$\lim_{t \rightarrow T_+} \frac{x(t)}{T_+ - t} = \mathbf{p}, \quad \lim_{t \rightarrow T_+} \frac{\lambda(t)}{T_+ - t} = 0,$$

and $(v_0, v_1) \in \mathcal{H} \cap C^\infty(\mathbb{R}^2 \setminus \{0\})$ with $(v_0 - u_\infty, v_1)$ compactly supported, such that

$$(i) \quad \inf \left\{ \left\| \vec{u}(t) - (v_0, v_1) - (Q_{\mathbf{p}}, \partial_t Q_{\mathbf{p}}) \right\|_{\mathcal{H}} : Q_{\mathbf{p}} \in \mathcal{M}_{\mathbf{p}} \right\} \xrightarrow{t \rightarrow T_+} 0,$$

$$(ii) \quad \left\| (\nabla u(t), \partial_t u(t)) - (\nabla v_0, v_1) \right\|_{L^2(\mathbb{R}^2 \setminus B_{\lambda(t)}(x(t)))} \xrightarrow{t \rightarrow T_+} 0,$$

where $B_{\lambda(t)}(x(t)) = \{x \in \mathbb{R}^2 : |x - x(t)| < \lambda(t)\}$, $\mathcal{M}_{\mathbf{p}}$ is the manifold of all transformations of $S_{\mathbf{p}}$ (by the group spanned by space translation, scaling, and S^2 isometries).