Semi-stabilization for the semilinear damped wave equation

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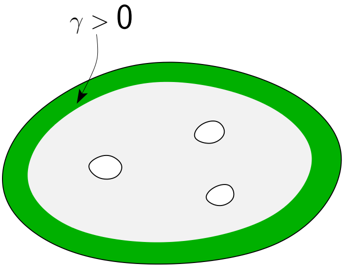
The semilinear damped wave equation

$$\boxed{\partial_{tt}^2 u + \gamma(x)\partial_t u = \Delta u - f(x,u)}$$



- Ω is a smooth compact manifold of dimension $d = 2$ with Dirichlet boundary conditions.
- the damping γ is in $\mathbb{L}^\infty(\Omega)$, $\gamma(x) \geq 0$
- f is smooth and of degree p

  $|f(x,u)| + |f'_x(x,u)| \leq C(1+|u|)^p$
  $|f'_u(x,u)| \leq C(1+|u|)^{p-1}$
- f is asymptotically of the sign of u:

  $\forall |u| \geq R \;, \quad f(x,u)u \geq 0$

Notations

$$\boxed{\partial^2_{tt}u + \gamma(x)\partial_t u = \Delta u - f(x,u)}$$

Set $X = H^1_0(\Omega) \times L^2(\Omega)$ and

$$U = \begin{pmatrix} u \\ \partial_t u \end{pmatrix} \qquad A = \begin{pmatrix} 0 & Id \\ \Delta & -\gamma(x) \end{pmatrix} \qquad F(U) = \begin{pmatrix} 0 \\ f(x,u) \end{pmatrix}$$

- $\Rightarrow$ e^{At} is a dissipative semigroup on X.
- $\Rightarrow$ Since f is of degree $p < \infty$ and Ω is of dimension $d = 2$, $F : X \longrightarrow X$ is defined and Lipschitz on the bounded sets.

We consider in X the equation

$$\boxed{\partial_t U = AU + F(U) \qquad U(t=0) = U_0 \in X}$$

The gradient dynamics

Set $V(x,u) = \int_0^u f(x,\xi)\,\mathrm{d}\xi$. The energy

$$\mathcal{E}(U) = \int_\Omega \frac{1}{2}(|\nabla u|^2 + |\partial_t u|^2) + V(x,u)\,\mathrm{d}x$$

is non-increasing along the trajectories since

$$\partial_t \mathcal{E}(U(t)) = -\int_\Omega \gamma(x)|\partial_t u|^2\,\mathrm{d}x$$

$\Rightarrow$ **Global existence of solutions**

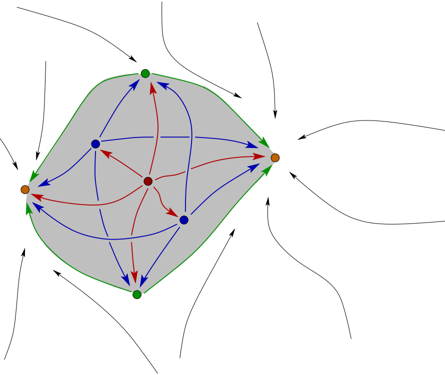
Motivations

The linear equation is **dissipative** and any solution goes to zero

$$\|e^{At}U\|_X \xrightarrow[t\longrightarrow+\infty]{} 0\ .$$

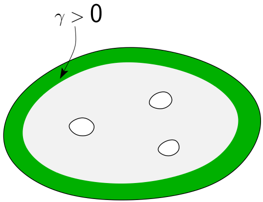
Do we still have stabilization of the nonlinear problem?

We expect gradient-like dynamics in X: there exists a compact global attractor and **any trajectory converges to the set of equilibrium points**.

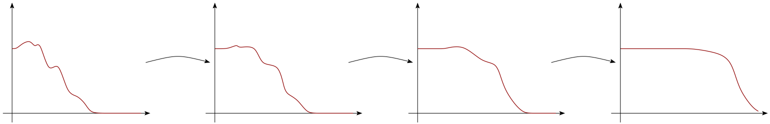


Motivations

Related problems:



- Control and stabilization problems.
- Does a linear dissipative (dispersive?) behaviour remains after addition of a nonlinearity?
- Convergence to equilibria/travelling fronts in nonlinear PDEs



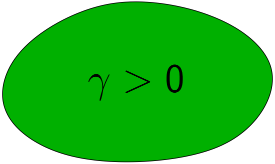
A historic result

A historic result

$$\boxed{\partial_{tt}^2 u + \gamma(x)\partial_t u = \Delta u - f(x,u)}$$

Assume:

- $\gamma(x) \geq \alpha > 0$ in Ω
- f is of degree $p < \infty$
- f is asymptotically of the sign of u



Theorem – J.K. Hale (1985) and A. Haraux (1985)

With the above assumptions, the dynamics of the damped wave equation admit a **compact global attractor** $\mathcal{A}$*. Moreover, it is gradient-like and* **any trajectory converges to the sets of equilibrium points***.*

A historic result

Step 1: the trajectories are bounded.

If f is asymptotically of the sign of u and of degree p, then the energy is well defined and bounded on bounded sets.

$$\frac{1}{2}\|U\|_X^2 + \min V \;\leq\; \int_\Omega \frac{1}{2}(|\nabla u|^2 + |\partial_t u|^2) + V(x,u)\,\mathrm{d}x \;\leq\; K(\|U\|_X)\ .$$

Since $\mathcal{E}$ is non-increasing, **the trajectories of bounded sets are bounded**.

A historic result

Step 2: the asymptotic compactness.

The linear semigroup is stabilized:

$$\forall t \geq 0 \ , \quad \|e^{At}\|_{\mathcal{L}(X)} \leq Me^{-\lambda t}$$

Moreover, if f is of degree p, then

$F : \begin{pmatrix} u \\ v \end{pmatrix} \in H_0^1(\Omega) \times L^2(\Omega) \longmapsto \begin{pmatrix} 0 \\ f(x,u) \end{pmatrix} \in H_0^1(\Omega) \times L^2(\Omega)$ is compact.

$$U(t) = e^{At}U_0 + \int_0^t e^{A(t-s)}F(U(s))\,\mathrm{d}s$$

$\Rightarrow$ **the bounded sets admits compact $\omega-$limit sets.**

A historic result

Step 3: a unique continuation property.

It is sufficient to show that the $\omega-$limit sets consists of equilibrium points. By Lasalle's principle, the trajectories $U(t) = (u, \partial_t u)$ in the $\omega-$limit sets have constant energy. So we have

$$\partial_t \mathcal{E}(U(t)) = -\int_\Omega \gamma(x)|\partial_t u|^2 \,\mathrm{d}x = 0 \ .$$

Since $\gamma(x) \geq \alpha > 0$, we have $\partial_t u \equiv 0$ and thus u is an equilibrium point.

A historic result

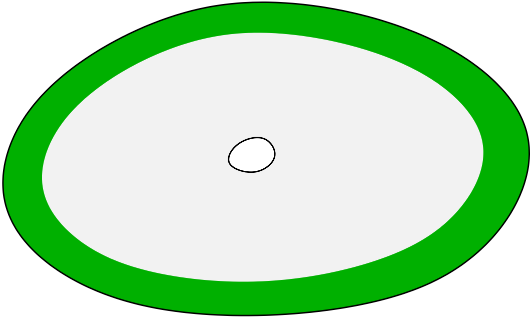
- Asymptotic compactness $\Leftrightarrow$ high frequencies are not really modified by the nonlinearity
- Unique continuation $\Leftrightarrow$ classify low-frequency solutions (equilibria, travelling fronts. . .)

Key arguments where we use γ positive:

1. $\|e^{At}\|_{\mathcal{L}(X)} \leq Me^{-\lambda t}$ has finite integral on $[0, +\infty)$
2. if $\mathcal{E}(U(t))$ is constant, then $\int \gamma(x)|u_t|^2 = 0$ and $u(t)$ is constant.

A standard extension

What happens when $\gamma(x)$ may vanish?



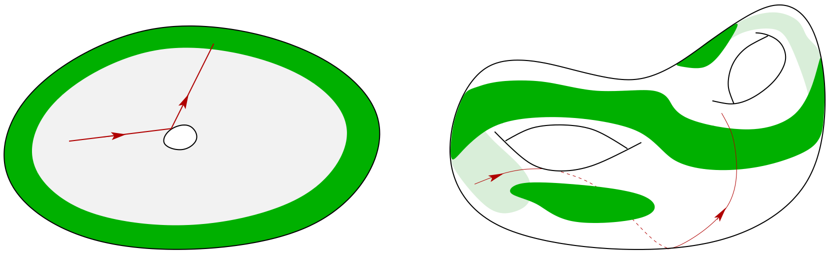
The decay of the linear semigroup

Theorem – J. Rauch and M. Taylor (1974)
C. Bardos, G. Lebeau and J. Rauch (1992)

$$\|e^{At}\|_{\mathcal{L}(X)} \leq Me^{-\lambda t}$$

$$\Longleftrightarrow$$

Any long enough geodesic meets the support of the damping γ



The unique continuation property

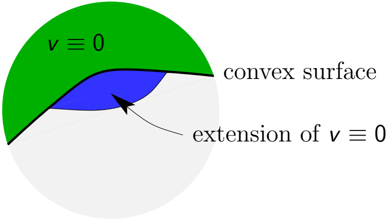
If $U(t)$ belongs to an $\omega-$limit set,

$$\partial_t \mathcal{E}(U(t)) = -\int_\Omega \gamma(x)|\partial_t u|^2 \,\mathrm{d}x = 0 \ .$$

So $v(t) = \partial_t u(t)$ vanishes in ω the support of γ. Thus, we have

$$v \equiv 0 \text{ in } \omega \times \mathbb{R} \qquad \text{and} \qquad \partial^2_{tt} v = \Delta v - f'_u(x, u(x,t))v \ .$$

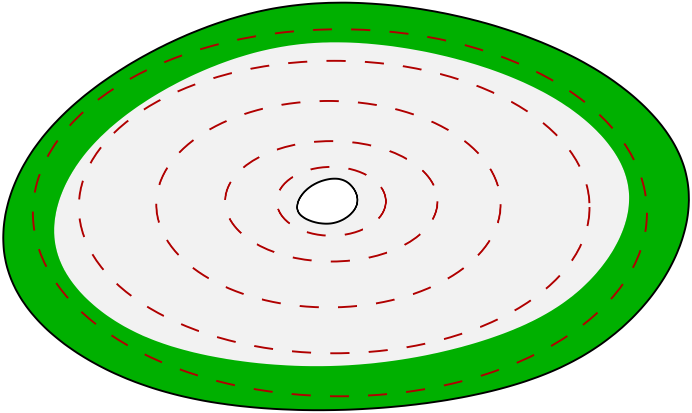
To conclude that $v \equiv 0$ everywhere, we need to use a unique continuation property.
Basically, we may extend the zone where $v \equiv 0$ through **convex surfaces**.



[N. Lerner and L. Robbiano, 1985], [L. Hörmander, 1985], [Tataru, 1996]

The unique continuation property

The gradient dynamics hold for the domain with zero or one hole

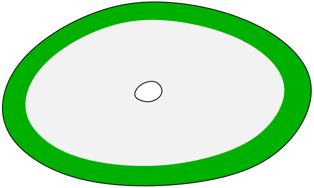


A classic result

$$\boxed{\partial^2_{tt} u + \gamma(x)\partial_t u = \Delta u - f(x, u)}$$

Assume:

- Ω is a two dimensional convex compact domain with or without a convex hole
- $\gamma(x) \geq \alpha > 0$ in a neighborhood of the exterior boundary of Ω.
- f of degree p
- f is asymptotically of the sign of u



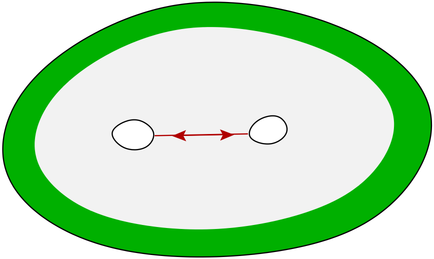
Theorem

With the above assumptions, the dynamics of the damped wave equation admit a **compact global attractor** $\mathcal{A}$*. Moreover, it is gradient-like and* **any trajectory converges to the sets of equilibrium points***.*

The disk with two holes

Without the geometric control condition

In some cases, the geometric control condition does not hold, but very few geodesics miss the support of the damping.



$$\|e^{At}U_0\|_{H^1\times L^2} \;\leq\; Me^{-\lambda\sqrt[3]{t}}\,\|U_0\|_{H^2\times H^1}$$

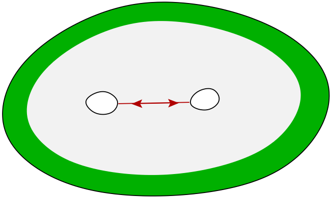
[N. Burq, 1993]
[N. Burq and M. Zworski, 2004]
[R.J. and C. Laurent, 2018]

The disk with two holes

$$\boxed{\partial^2_{tt} u + \gamma(x)\partial_t u = \Delta u - f(x,u)}$$

Assume:

- Ω is a convex compact domain of dimension 2 with two convex holes
- $\gamma(x) \geq \alpha > 0$ in a neighborhood of the exterior boundary of Ω.
- f of degree $p < \infty$
- f is asymptotically of the sign of u



Theorem – R.J. and C. Laurent (2018)

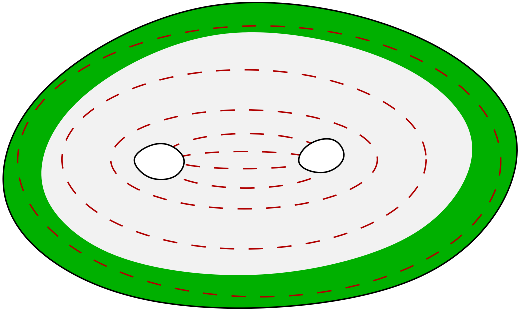
With the above assumptions, the dynamics of the damped wave equation are gradient-like. Moreover, they admit a compact global $(D(A), X)$−attractor $\mathcal{A}$ in the following sense: any bounded set of $D(A) = (H^2 \cap H^1_0) \times H^1_0$ is attracted by $\mathcal{A}$ in the norm $X = H^1_0 \times L^2$. Moreover, **any trajectory converges to the sets of equilibrium points** *and the convergence is uniform in balls of $D(A)$.*

The disk with two holes

Main arguments:

$$\|e^{At}U_0\|_{H^1\times L^2} \leq Me^{-\lambda t^{1/3}}\|U_0\|_{H^2\times H^1}$$

$$U(t) = e^{At}U_0 + \int_0^t e^{A(t-s)}F(U(s))\,\mathrm{d}s$$



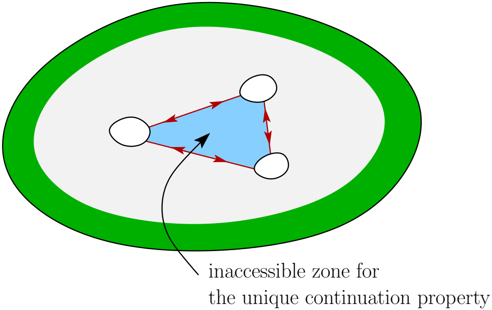
$+$ some technical tricks

The disk with three holes

If there are three holes or more, with additional technical assumptions, we still have the decay

$$\|e^{At}U_0\|_{H^1\times L^2} \;\leq\; Me^{-\lambda\sqrt[3]{t}}\,\|U_0\|_{H^2\times H^1}$$

But the unique continuation property of Lerner-Robbiano-Hörmander-Tataru cannot be used:



An analytic unique continuation property

Theorem – L. Robbiano and C. Zuily (1998) L. Hörmander (1997)

Assume $\omega \neq \emptyset$ *and* $v(t) = \partial_t u(t)$ *solves*

$$v \equiv 0 \text{ in } \omega \times \mathbb{R} \quad \text{ and } \quad \partial^2_{tt} v = \Delta v - f'_u(x, u(x,t))v \ .$$

Assume moreover that $t \mapsto f'_u(x, u(x,t))$ *is analytic then* $v \equiv 0$ *everywhere.*

[J.K. Hale and G. Raugel, 2003] let us hope that if $f(x, u)$ is analytic in u, then a function u in the attractor should be analytic in time and thus $f'_u(x, u(x,t))$ is also analytic.

An analytic unique continuation property

In the proofs of [J.K. Hale and G. Raugel, 2003], a global solution u is split between the **low-frequencies** $P_n u$ **and the high-frequencies** $Q_n u$. It is used that

$$\|e^{At}U\|_X \leq Me^{-\lambda t}\|U\|_X \implies \|e^{Q_nAQ_nt}Q_nU\|_X \leq Ne^{-\mu t}\|Q_nU\|_X \ .$$

In our case, we would like to obtain

$$\|e^{At}U\|_X \leq Me^{-\lambda \sqrt[3]{t}}\|U\|_{D(A)} \implies \|e^{Q_nAQ_nt}Q_nU\|_X \leq h(t)\|Q_nU\|_{D(A)} \ .$$

$\implies$ we adapt the ideas of [J.K. Hale and G. Raugel, 2003] but several technical problems have to be overcome.

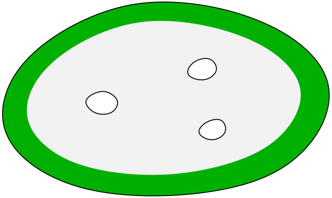
[C.J.K. Batty and Th. Duyckaerts, 2008], [A. Borichev and Y. Tomilov, 2010], [N. Anantharaman and M. Léautaud, 2014]

The disk with three holes

$$\boxed{\partial_{tt}^2 u + \gamma(x)\partial_t u = \Delta u - f(x,u)}$$

Assume:

- Ω is as opposite and the holes are not aligned and small enough
- $f(x,u)$ is analytic in u
- f of degree $p < \infty$
- f is asymptotically of the sign of u



Theorem – R.J. and C. Laurent (2018)

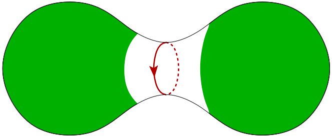
With the above assumptions, the dynamics admit a compact global $(D(A), X)-$attractor. Moreover, **any trajectory converges to the sets of equilibrium points** *and the convergence is uniform in balls of $D(A)$.*

Conclusion

Additional results

- **Other geometries are possible**

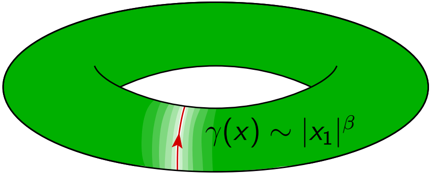
The peanut of rotation



$\|e^{At}U_0\|_{H^1\times L^2} \leq Me^{-\lambda\sqrt{t}}\|U_0\|_{H^2\times H^1}$

[E. Schenck, 2011]
[H. Christianson, E. Schenck, A. Vasy and J. Wunsch, 2014]

The torus with degenerated damping



$\|e^{At}U_0\|_{H^1\times L^2} \leq \frac{C}{(1+t)^{1+2/\beta}}\|U_0\|_{H^2\times H^1}$

[M. Léautaud and N. Lerner, 2015]

- **Higher dimension**

  In dimension $d = 3$, assume that f is Sobolev-subcritical, that is of degree p with $p < 3$. It should also be possible to go to f energy-subcritical, that is of degree p with $p < 5$ by using Strichartz estimates, see [B. Dehman, G. Lebeau and E. Zuazua, 2003], [R.J. and C. Laurent, 2013]

Open problem

$$U(t) = e^{At}U_0 + \int_0^t e^{A(t-s)}F(U(s))\,\mathrm{d}s$$

Open question n.1:

How important is the integrability of the linear decay?

For example, if the linear decay is simply

$$\|e^{At}U_0\|_{H^1\times L^2} \leq \frac{C}{\ln(2+t)}\,\|U_0\|_{H^2\times H^1}$$

does the asymptotic compactness hold?

Open problem

Open question n.2:

In dispersive equations or if F not compact, can we use appropriate estimates to adapt the proofs?

Thanks for your attention!

- R.J. and C. Laurent, *Semi-uniform decay for some semilinear damped wave equations*, almost preprint.
- R.J. and C. Laurent, *A note on the global controllability of the semilinear wave equation*, SIAM Journal on Control and Optimization n°52 (2014), pp. 439–450.
- R.J. and C. Laurent, *Stabilization for the semilinear wave equation with geometric control condition*, Analysis and PDE n°6 (2013), pp. 1089–1119.